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Solutions 191 Problems of Chapter 3

3.1 Introducing a Stokes damping into Newton ' s second law (3.1) leads to the real part of the refractive index

$n(\omega) = \text{Re} \left[1 + \chi^{(1)}(\omega) \right]$, with the susceptibility $\chi^{(1)}(\omega) = \frac{e^2 N_0}{m^2 \epsilon_0 \omega^2} \left[\frac{1}{\omega_0 - \omega - i\gamma} + \frac{1}{\omega_0 + \omega - i\gamma} \right]$ (compare (3.3)). The absorption coefficient is

$\alpha(\omega) = 2 \frac{e^2 N_0}{m^2 \epsilon_0 \omega^2} \frac{\gamma}{(\omega_0 - \omega)^2 + \gamma^2}$ (compare (3.3)). The absorption coefficient is

$\alpha(\omega) = 2 \frac{e^2 N_0}{m^2 \epsilon_0 \omega^2} \frac{\gamma}{(\omega_0 - \omega)^2 + \gamma^2}$. As usual,

phase and group velocities are given

by $v_{\text{phase}}(\omega) = \frac{\omega}{k} = \dots$

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Manual - Maharashtra

Nonlinear optical phenomena are “ nonlinear ” in the sense that they occur when the response of a material system to an applied optical field depends in a nonlinear manner on the strength of the optical field. In

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Nonlinear optics, the optical response can be described by expressing the polarization as a power series in the field strength.

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